Using covering spaces to model soap films

Maurizio Paolini (maurizio.paolini@unicatt.it)

Dipartimento di Matematica e Fisica "Niccolò Tartaglia" Università Cattolica del Sacro Cuore, Brescia

Okinawa, march 7, 2017

- 4 同 ト 4 ヨ ト 4 ヨ ト

Working in our group:

S. Amato, G. Bellettini, M. P., Constrained BV functions on covering spaces for minimal networks and Plateau's type problems, Adv. Calc. Var.

Marcello Carioni, Franco Pasquarelli, Alessandra Pluda

・ 同 ト ・ ヨ ト ・ ヨ ト

Working in our group:

S. Amato, G. Bellettini, M. P., *Constrained BV functions on covering spaces for minimal networks and Plateau's type problems*, Adv. Calc. Var. Marcello Carioni, Franco Pasguarelli, Alessandra Pluda

Working on the same subject:

K. Brakke, *Soap films and covering spaces*, Journal of Geometric Analysis, 1995

伺 ト イヨト イヨト

Working in our group:

S. Amato, G. Bellettini, M. P., *Constrained BV functions on covering spaces for minimal networks and Plateau's type problems*, Adv. Calc. Var. Marcello Carioni, Franco Pasquarelli, Alessandra Pluda

Working on the same subject:

K. Brakke, *Soap films and covering spaces*, Journal of Geometric Analysis, 1995

A. Chambolle, D. Cremers, T. Pock, *A convex approach to minimal partitions*, SIAM Journal on Imaging Sciences, 2012

高 とう モン・ く ヨ と

- Motivation
- Covering space
- Knots and links
- The tripod (model for the triple junction)
- The convexification problem
- Other examples

- 4 回 2 - 4 □ 2 - 4 □

Find a material interface (say soap film) with minimal area and given boundary in 3D Corresponding evolution problem (Mean Curvature Flow) Using the BV machinery [De Giorgi], suitable e.g. for relaxation via diffused interface **Features:** Knotted curves in 3D, triple junctions

・ 同 ト ・ ヨ ト ・ ヨ ト …

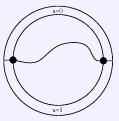
Find a material interface (say soap film) with minimal area and given boundary in 3D

Corresponding evolution problem (Mean Curvature Flow)

Using the BV machinery [De Giorgi], suitable e.g. for relaxation via diffused interface

Features: Knotted curves in 3D, triple junctions

A "too" simple example



 Γ is the equator of the unit sphere Ω is an enlarged sphere Force u = 1 and u = 0 outside the unit sphere Minimize the total variation in $BV(\Omega; \{0, 1\})$

subject to the constraint

The jump set $\boldsymbol{\Sigma}$ of a minimizer is the desired soap film

(4月) イヨト イヨト

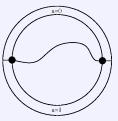
Find a material interface (say soap film) with minimal area and given boundary in 3D

Corresponding evolution problem (Mean Curvature Flow)

Using the BV machinery [De Giorgi], suitable e.g. for relaxation via diffused interface

Features: Knotted curves in 3D, triple junctions

A "too" simple example



 Γ is the equator of the unit sphere

 Ω is an enlarged sphere

Force u = 1 and u = 0 outside the unit sphere

Minimize the total variation in $BV(\Omega; \{0, 1\})$

subject to the constraint

The jump set Σ of a minimizer is the desired soap film

イロト イヨト イヨト イヨト

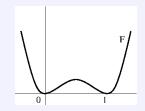
We trick the Plateau problem into a phase-separation problem, the two phases being $\{u = 0\}$ and $\{u = 1\}$

Allen-Cahn equation: brief review

Reaction/diffusion equation arising in the context of **phase transitions** with a diffused interface:

 $\begin{cases} \epsilon \partial_t u - \epsilon \Delta u + \frac{1}{\epsilon} f(u) = 0 & \text{in } \Omega \\ + \text{ initial and boundary conditions} \end{cases}$

- *u*: order parameter (phase indicator),
- Ω : domain in \mathbb{R}^d , d = 2, 3,
- $\epsilon > 0$: small relaxation parameter,
- f = F': derivative of a double equal well potential F (or double-obstacle: deep quench limit [Elliott et al]).



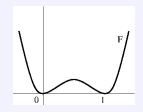
The solution u exhibits a thin transition layer $\mathcal{O}(\epsilon)$ -wide between the phase upprox 0 and the phase upprox 1

Allen-Cahn equation: brief review

Reaction/diffusion equation arising in the context of **phase transitions** with a diffused interface:

 $\begin{cases} \epsilon \partial_t u - \epsilon \Delta u + \frac{1}{\epsilon} f(u) = 0 & \text{in } \Omega \\ + \text{ initial and boundary conditions} \end{cases}$

- *u*: order parameter (phase indicator),
- Ω : domain in \mathbb{R}^d , d = 2, 3,
- $\epsilon > 0$: small relaxation parameter,
- f = F': derivative of a double equal well potential F (or double-obstacle: deep quench limit [Elliott et al]).



イロト イポト イヨト イヨト

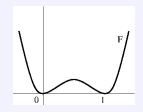
The solution u exhibits a thin transition layer $\mathcal{O}(\epsilon)$ -wide between the phase upprox 0 and the phase upprox 1

Allen-Cahn equation: brief review

Reaction/diffusion equation arising in the context of **phase transitions** with a diffused interface:

 $\begin{cases} \epsilon \partial_t u - \epsilon \Delta u + \frac{1}{\epsilon} f(u) = 0 & \text{in } \Omega \\ + \text{ initial and boundary conditions} \end{cases}$

- *u*: order parameter (phase indicator),
- Ω : domain in \mathbb{R}^d , d = 2, 3,
- $\epsilon > 0$: small relaxation parameter,
- f = F': derivative of a double equal well potential F (or double-obstacle: deep quench limit [Elliott et al]).



イロト イポト イヨト イヨト

The solution u exhibits a thin transition layer $\mathcal{O}(\epsilon)$ -wide between the phase $u \approx 0$ and the phase $u \approx 1$

Singular limit $\epsilon \rightarrow 0$

The transition layer approximates a sharp interface that moves by mean curvature:

$$V = -\kappa$$

[X. Chen, Bronsard-Kohn, Evans-Soner-Souganidis,] [Barles-Soner-Souganidis, ...]

Optimal $\mathcal{O}(\epsilon^2)$ or quasi-optimal $\mathcal{O}(\epsilon^2 | \log \epsilon |)$ error estimate. [Nochetto-P.-Verdi, Nochetto-Verdi, Bellettini-P.]

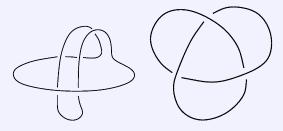
A.C. is the gradient flow of

 $\mathcal{F}_{\epsilon}(u) := \tfrac{\epsilon}{2} \int_{\Omega} |\nabla u|^2 \, dx + \tfrac{1}{\epsilon} \int_{\Omega} F(u) \, dx$

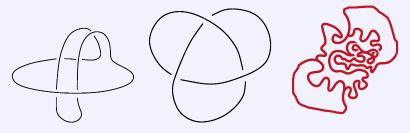
Convergence as $\epsilon \rightarrow 0$

 \mathcal{F}_{ϵ} Γ -converges to $c_0 \mathcal{F}$ with $\mathcal{F}(u) := \int_{\Omega} |Du|$, c_0 a suitable constant depending on F, $u \in BV(\Omega, \{0, 1\})$

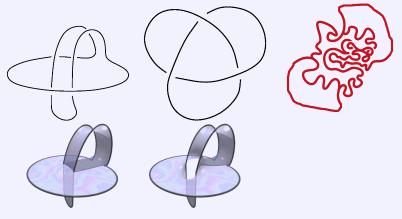
This simple BV approach for the circle example does not work in general. The curve Γ must be such that it lies in the boundary of a convex (or at least mean-convex) body in \mathbb{R}^3



This simple BV approach for the circle example does not work in general. The curve Γ must be such that it lies in the boundary of a convex (or at least mean-convex) body in \mathbb{R}^3



This simple BV approach for the circle example does not work in general. The curve Γ must be such that it lies in the boundary of a convex (or at least mean-convex) body in \mathbb{R}^3



[images courtesy of Emanuele Paolini]

Problem: the desired surface Σ does not naturally separate two phases near Γ

Idea

Work on a "covering space" Y of $\Omega = \mathbb{R}^3 \setminus \Gamma$ (if regularity of $\partial \Omega$ is needed Γ can be replaced by a tubular neighborhood Γ_{ϵ})

 $\Omega = B \setminus \Gamma_{\epsilon}$, where B is a large ball containing Γ_{ϵ} or $B = \mathbb{S}^3$ $p: Y \to \Omega$ is a covering (of finite degree k)

V

Find $u \in BV(Y; \{0,1\})$ such that

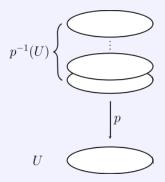
- u is given near the boundary of Ω
- *u* satisfies the constraint

$$\sum_{x \in p^{-1}(x)} u(y) = 1$$

on the fiber above any $x\in \Omega$

 Ω is path-connected. $p: Y \rightarrow \Omega$ is locally trivial

For any $x \in \Omega$ there is a small neighborhood U such that $p^{-1}(U)$ is topologically the disjoint union of kcopies of U. k (degree of the covering) is locally constant in Ω , hence constant. It is not finite in general, but for our purposes $k < \infty$



A covering can be globally trivial: Y is the disjoint union of k copies of Ω , but the interesting case is when Y is connected.

イロト イヨト イヨト イヨト

A covering can be globally trivial: Y is the disjoint union of k copies of Ω , but the interesting case is when Y is connected. **Example** (k = 2) $\Omega = \mathbb{S}^2 \setminus \{A, B\}$

- Start with two copies of Ω (deck 1 and deck 2)
- cut along the segment AB
- glue deck 1 above with deck 2 below and viceversa

We obtain a nontrivial covering of degree 2



伺 ト イヨト イヨト



Maurizio Paolini (maurizio.paolini@unicatt.it) Using covering spaces to model soap films

Important property If $\gamma : [0,1] \to \Omega$ is a closed curve with $\gamma(0) = \gamma(1) = x$, then after chosing $y \in p^{-1}(x)$ we can "lift" γ into $\gamma_{\#} : [0,1] \to Y$ with $\gamma_{\#}(0) = y$, continuous and such that $p \circ \gamma_{\#} = \gamma$. In general $\gamma_{\#}(1) \in p^{-1}(x)$ is different from $\gamma_{\#}(0)$ and independent on continuous deformations of γ : $\gamma \in \pi_1(\Omega, x)$ acts on the fiber at x

 $\pi_1(\Omega, x)$ is the fundamental group of Ω with base point x

通 とう ほう うちょう

Important property If $\gamma : [0,1] \to \Omega$ is a closed curve with $\gamma(0) = \gamma(1) = x$, then after chosing $y \in p^{-1}(x)$ we can "lift" γ into $\gamma_{\#} : [0,1] \to Y$ with $\gamma_{\#}(0) = y$, continuous and such that $p \circ \gamma_{\#} = \gamma$. In general $\gamma_{\#}(1) \in p^{-1}(x)$ is different from $\gamma_{\#}(0)$ and independent on continuous deformations of γ : $\gamma \in \pi_1(\Omega, x)$ acts on the fiber at $x = \pi_1(\Omega, x)$ is the fundamental group of Ω with base point x

Trivial example

If Ω is simply connected \Rightarrow covering is globally trivial

イロト イポト イヨト イヨト

Basic example

Remark

In this perspective the "point at infinity" of \mathbb{R}^2 is important. We prefer to work in \mathbb{S}^2 . Conversely using \mathbb{R}^3 or \mathbb{S}^3 is basically equivalent (in dimension 3 a curve cannot be obstructed by a single point)

向下 イヨト イヨト

Basic example

Remark

In this perspective the "point at infinity" of \mathbb{R}^2 is important. We prefer to work in \mathbb{S}^2 . Conversely using \mathbb{R}^3 or \mathbb{S}^3 is basically equivalent (in dimension 3 a curve cannot be obstructed by a single point)

Knot/link K, k = 2

- Take a Seifert surface S of the knot
- Start with two copies of $\Omega := \mathbb{R}^3 \setminus K$ (deck 1 and deck 2)
- Cut them along S and glue back with a deck exchange

- 4 同 6 4 日 6 4 日 6

Basic example

Remark

In this perspective the "point at infinity" of \mathbb{R}^2 is important. We prefer to work in \mathbb{S}^2 . Conversely using \mathbb{R}^3 or \mathbb{S}^3 is basically equivalent (in dimension 3 a curve cannot be obstructed by a single point)

Knot/link K, k = 2

- Take a Seifert surface S of the knot
- Start with two copies of $\Omega := \mathbb{R}^3 \setminus K$ (deck 1 and deck 2)
- Cut them along S and glue back with a deck exchange

Equivalently: Start with $X = \{ \text{paths } \gamma \text{ starting at } x \}$

- $\gamma_1 \sim \gamma_2$ if $\gamma_1(1) = \gamma_2(1)$ and
- $\gamma_1\gamma_2^{-1}$ has **even** linking number with K

Define
$$Y=X/\sim$$
, $p(\gamma)=\gamma(1)\in \Omega$

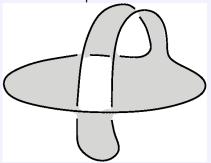


In this case γ has even linking number iff the corresponding word in the group presentation has even length

< a, b; aba = bab > is a Wirtinger presentation of the trefoil knot; < x, y; $x^2 = y^3$ > is **not** a Wirtinger p. of the trefoil knot

向下 イヨト イヨト

There are many ways of constructing a Seifert surface, one is shown in the picture



Now that we have a "double" (k = 2) covering of Ω we shall consider functions $u \in S_0$ with

$$S_0 := \left\{ u \in BV(Y; \{0,1\}) : \sum_{y \in p^{-1}(x)} u(y) = 1 \right\}$$

Functions in S_0 **must** jump when circling Γ once along small curves \Rightarrow the jump set must touch Γ at all points. We have no control on the topology, but this is not a problem

伺 ト イヨト イヨト

The BV setting 2

For $u \in S_0$ we can define the energy

$$\mathcal{F}(u) := rac{1}{2} \int_{Y} |Du|$$

Due to the constraint the jump set projects nicely on Ω and we account the projected jump twice, hence the factor 1/2The coarea formula allows to convexify S_0 into

$$S:=\left\{u\in BV(Y,[0,1]):\sum_{p(y)=x}u(y)=1
ight\}$$

(same energy) and get an essentially equivalent problem

Warning!

This is specific to the
$$k = 2$$
 case

イロト イポト イヨト イヨト

Strictly related: he works in the context of currents on the covering space Y, the constraint is imposed by constructing a higher-level covering W of Y made of "pairs" of distinct points on each fiber This approach shows its power in the case k > 2 but makes sense also for k = 2

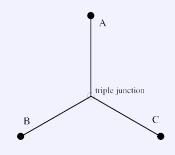
The boundary of a flat chain in Y can be associated (in infinite ways) to a current T in W and we seek to minimize the mass of T (that represents the soap film itself) obtaining the "film mass"

向下 イヨト イヨト

Model problem to tackle triple junctions:

 $\Omega := \mathbb{S}^2 \setminus \{A, B, C\}$

where A, B, C are the vertices of an equilateral triangle. The solution that we would like to see is the "tripod" To guarantee that Γ touches all three points coverings of degree 2 do not work...



向下 イヨト イヨト

This is the solution of the Steiner problem (angles of 120 degrees)

A covering for the tripod

For k = 3 a natural construction by cut & paste is obtained by cutting three copies of \mathbb{S}^3 along two of the three sides, say AB and BC and glue them again according to given permutations σ_1 and σ_2 of the three strata

A B C n₂

向下 イヨト イヨト

The only working choice turns out to be

$$\sigma_1 = (1, 2, 3), \qquad \sigma_2 = \sigma_1^{-1} = (1, 3, 2)$$

and again minimize $\mathcal{F}(u) := rac{1}{2} \int_{Y} |Du|$ on the domain

$$S_0 := \left\{ u \in BV(Y; \{0,1\}) : \sum_{p(y)=x} u(y) = 1 \right\}$$

[show video]

Convexification is tricky

$$S:=\left\{u\in BV(Y;[0,1]):\sum_{p(y)=x}u(y)=1
ight\}$$

If we minimize $\int_{Y} |Du|$, the possibility of mixing all three "phases" allows to obtain a lower energy than for the nonconvex problem

向下 イヨト イヨト

Convexification is tricky

$$S:=\left\{u\in BV(Y;[0,1]):\sum_{p(y)=x}u(y)=1
ight\}$$

If we minimize $\int_{Y} |Du|$, the possibility of mixing all three "phases" allows to obtain a lower energy than for the nonconvex problem

Brakke: "film mass"

・ 同 ト ・ ヨ ト ・ ヨ ト

Convexification is tricky

$$S:=\left\{u\in BV(Y;[0,1]):\sum_{p(y)=x}u(y)=1
ight\}$$

If we minimize $\int_{Y} |Du|$, the possibility of mixing all three "phases" allows to obtain a lower energy than for the nonconvex problem

Brakke: "film mass"

Chambolle et al.: "local convex envelope"

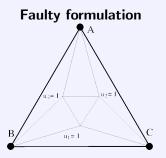
 \Rightarrow the same energy (extension of \mathcal{F} to the convexified domain S) in the case of the tripod

(1) マン・ション・

Trivializing the tripod example

Since this example is very special, we can restate it in a more usual setting by restricting Ω to Ω_T , the triangle *ABC*, thus obtaining a globally trivial 3-covering with restrictions of u on the three decks that we can denote

$$u_1, u_2, u_3 \in BV(\Omega_T, [0, 1]), \qquad \sum_i u_i = 1$$



Mixture of all three phases inside the small triangle leads to total variation smaller than that of the tripod

Identifying the convexified

The Brakke's "film mass" turns out to be

$$\mathcal{F}(u) = sup_{\varphi \in K} \sum_{i} \int_{\Omega_{\mathcal{T}}} u_{i} \operatorname{div} \varphi_{i} dx$$

where

$$\mathcal{K} = \{\varphi_i \in [\mathcal{C}_c^{\infty}(\Omega_{\mathcal{T}})]^3, i = 1, 2, 3 : |\varphi_i(x) - \varphi_j(x)|_2 \le 2 \,\forall i \neq j, x \in \Omega\}$$

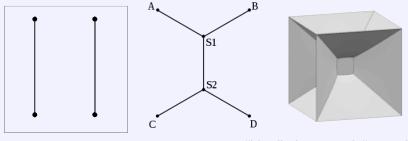
In the case $u_1 \in W^{1,1}$ a lengthy computation shows that this can be computed as

$$\mathcal{F}(u) = \frac{1}{3} \int_{\Omega_T} \text{steinerdist}(\nabla(u_1 - u_2), \nabla(u_2 - u_3), \nabla(u_3 - u_1)) \, dx$$

consistent with the so-called local convex envelope of Chambolle et al.



Other examples



"blow" the central "square"

▲圖> ▲屋> ▲屋>

Э

THANK YOU FOR YOUR PATIENCE!

・ロン ・四 と ・ 回 と ・ 回 と

3

Maurizio Paolini (maurizio.paolini@unicatt.it) Using covering spaces to model soap films

◆□▶ ◆□▶ ◆三▶ ◆三▶ 三三 - のへで